## Differential Equations

 Class Notes

Auxiliary Equations with Complex Roots (Section 4.3)
In the previous section, we put our mass-spring oscillators on hold to study only exponential solutions to the linear, second-order constant coefficient equations. Here, we see that these massspring systems can give rise to an auxiliary equation that has complex roots.

## Rationale for Solutions:

For the linear, second-order diff. eq. $a y^{\prime \prime}+b y^{\prime}+c y=0$, we have its auxiliary equation $a r^{2}+b r+c=0$. When $b^{2}-4 a c<0$, the equation has two complex roots
$r=\frac{-b \pm \sqrt{b^{2}-4 a c}}{2 a}=\frac{-b}{2 a} \pm \frac{\sqrt{4 a c-b^{2}}}{2 a} \cdot i$.
We will call these roots $r_{1}=\alpha+i \beta$ and $r_{2}=\alpha-i \beta$, defining $\alpha=-b / 2 a$ and $\beta=\frac{\sqrt{4 a c-b^{2}}}{2 a}$.
By these definitions, we see that $\alpha$ and $\beta$ are real numbers.

From the previous section, the solutions to our diff. eq. are $y_{1}(t)=e^{r_{1} t}$ and $y_{2}(t)=e^{r_{2} t}$. Using the values we have developed for our roots, these solutions are now in the form $y_{1}(t)=e^{(\alpha+i \beta) t}$ and $y_{2}(t)=e^{(\alpha-i \beta) t}$.

We can use the Maclaurin series and Euler's formula (described in the book) to further rewrite our solutions. We now have $y_{1}(t)=e^{(\alpha+i \beta) t}=e^{\alpha t}(\cos (\beta t)+i \sin (\beta t))$ and $y_{2}(t)=e^{(\alpha-i \beta) t}=e^{\alpha t}(\cos (\beta t)-i \sin (\beta t))$.
$\circ \bigcirc$


These solutions involve complex numbers. Do they have to?

Finishing this out, we can say that $y(t)=c_{1} \cdot y_{1}(t)+c_{2} \cdot y_{2}(t)$ is a general solution to the diff. eq. $a y^{\prime \prime}+b y^{\prime}+c y=0$.

As the thought bubble at the bottom of the previous page ponders, we would like to find realvalued functions that are solutions to our diff. eq.. In fact, we have this lemma.

## Lemma 2: Real Solutions Derived from Complex Solutions:

Let $z(t)=u(t)+i \cdot v(t)$ be a solution to $a y^{\prime \prime}+b y^{\prime}+c y=0$, where $a$, $b$, and $c$ are real numbers. Then the real part (which is $u(t)$ ) and the imaginary part (which is $v(t)$ ) are real-valued solutions to the diff. eq..

This leads to our main theorem.


## Theorem: Complex Conjugate Roots:

If the auxiliary equation has complex roots $\alpha \pm i \beta$, then two linearly independent solutions to $a y^{\prime \prime}+b y^{\prime}+c y=0$ are $y_{1}(t)=e^{\alpha t} \cos (\beta t)$ and $y_{2}(t)=e^{\alpha t} \sin (\beta t)$.

A general solution is given by $y(t)=c_{1} e^{\alpha t} \cos (\beta t)+c_{2} e^{\alpha t} \sin (\beta t)$ where $c_{1}$ and $c_{2}$ are real constants.
expl 1: The auxiliary equation for this diff. eq. has complex roots. Find a general solution. $y^{\prime \prime}-4 y^{\prime}+7 y=0$
$\operatorname{expl} 2$ : Solve the initial value problem.

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y^{\prime \prime}+9 y=0, \quad y(0)=1, \quad y^{\prime}(0)=1
$$



## Vibrating Springs without Damping:

expl 3: A vibrating spring without damping can be modeled by the diff. eq. $m y^{\prime \prime}+b y^{\prime}+k y=0$. By taking $b=0$ because there is no damping, this equation becomes $m y^{\prime \prime}+k y=0$.
a.) If $m=10 \mathrm{~kg}, k=250 \mathrm{~kg} / \mathrm{sec}^{2}, y(0)=0.3 \mathrm{~m}$, and $y^{\prime}(0)=-0.1 \mathrm{~m} / \mathrm{sec}$, find the equation of motion for this undamped vibrating spring.
b.) After how many seconds will the mass in part $a$ first cross the equilibrium point?
c.) When the equation of motion is of the form $y(t)=c_{1} e^{\alpha t} \cos (\beta t)+c_{2} e^{\alpha t} \sin (\beta t)$, the motion is said to be oscillatory with frequency $\beta / 2 \pi$. Find the frequency of oscillation
 for the spring system in part $a$.

